## The Fourier Transform

- Introduction
- Orthonormal bases for $\mathbb{R}^{n}$
- Inner product
- Length
- Orthogonality
- Change of basis
- Matrix transpose
- Complex vectors
- Orthonormal bases for $\mathbb{C}^{n}$
- Inner product
- Hermitian transpose
- Orthonormal bases for $2 \pi$ periodic functions
- Shah basis
- Harmonic signal basis
- Fourier series
- Fourier transform


## Orthonormal bases for $\mathbb{R}^{n}$

Let $\mathbf{u}=\left[u_{1}, u_{2}\right]^{\mathrm{T}}$ and $\mathbf{v}=\left[v_{1}, v_{2}\right]^{\mathrm{T}}$ be vectors in $\mathbb{R}^{2}$. We define the inner product of $\mathbf{u}$ and $\mathbf{v}$ to be

$$
\langle\mathbf{u}, \mathbf{v}\rangle=u_{1} v_{1}+u_{2} v_{2} .
$$

We can use the inner product to define notions of length and angle. The length of $\mathbf{u}$ is given by the square root of the inner product of $\mathbf{u}$ with itself:

$$
\begin{aligned}
|\mathbf{u}| & =\langle\mathbf{u}, \mathbf{u}\rangle^{\frac{1}{2}} \\
& =\sqrt{u_{1}^{2}+u_{2}^{2}}
\end{aligned}
$$

The angle between $\mathbf{u}$ and $\mathbf{v}$ can also be defined in terms of inner product:

$$
\langle\mathbf{u}, \mathbf{v}\rangle=|\mathbf{u}||\mathbf{v}| \cos \theta
$$

where

$$
\theta=\cos ^{-1}\left(\frac{\langle\mathbf{u}, \mathbf{v}\rangle}{|\mathbf{u}||\mathbf{v}|}\right) .
$$

## Orthogonality

An important special case occurs when

$$
\langle\mathbf{u}, \mathbf{v}\rangle=|\mathbf{u}||\mathbf{v}| \cos \theta=0
$$

When $\cos \theta$ equals zero, $\theta=\pi / 2=90^{\circ}$.

## Orthonormal bases for $\mathbb{R}^{n}$

Any $n$ orthogonal vectors which are of unit length

$$
\left\langle\mathbf{u}_{i}, \mathbf{u}_{j}\right\rangle=\left\{\begin{array}{l}
1 \text { if } i=j \\
0 \text { otherwise } .
\end{array}\right.
$$

form an orthonormal basis for $\mathbb{R}^{n}$. Any vector in $\mathbb{R}^{n}$ can be expressed as a weighted sum of $\mathbf{u}_{1}$, $\mathbf{u}_{2}, \mathbf{u}_{3}, \ldots, \mathbf{u}_{n}$ :

$$
\mathbf{v}=w_{1} \mathbf{u}_{1}+w_{2} \mathbf{u}_{2}+w_{3} \mathbf{u}_{3}+\ldots+w_{n} \mathbf{u}_{n} .
$$

- Question How do we find $w_{1}, w_{2}, w_{3}, \ldots, w_{n}$ ?
- Answer Using inner product.


## Example

Consider two orthonormal bases. The first basis is defined by the vectors $\mathbf{u}_{1}=\left[\begin{array}{l}1 \\ 0\end{array}\right]$ and $\mathbf{u}_{2}=$ $\left[\begin{array}{l}0 \\ 1\end{array}\right]$. It is easy to verify that these two vectors form an orthonormal basis:

$$
\begin{aligned}
& \left\langle\left[\begin{array}{l}
1 \\
0
\end{array}\right],\left[\begin{array}{l}
0 \\
1
\end{array}\right]\right\rangle=1 \cdot 0+0 \cdot 1=0 \\
& \left\langle\left[\begin{array}{l}
1 \\
0
\end{array}\right],\left[\begin{array}{l}
1 \\
0
\end{array}\right]\right\rangle=1 \cdot 1+0 \cdot 0=1 \\
& \left\langle\left[\begin{array}{l}
0 \\
1
\end{array}\right],\left[\begin{array}{l}
0 \\
1
\end{array}\right]\right\rangle=0 \cdot 0+1 \cdot 1=1 .
\end{aligned}
$$

## Example (contd.)

The second, by the vectors $\mathbf{u}_{1}^{\prime}=\left[\begin{array}{c}\cos \theta \\ \sin \theta\end{array}\right]$ and $\mathbf{u}_{2}^{\prime}=\left[\begin{array}{r}-\sin \theta \\ \cos \theta\end{array}\right]$. It is also easy to verify that these two vectors form an orthonormal basis:
$\left\langle\left[\begin{array}{c}\cos \theta \\ \sin \theta\end{array}\right],\left[\begin{array}{r}-\sin \theta \\ \cos \theta\end{array}\right]\right\rangle=-\cos \theta \sin \theta+\cos \theta \sin \theta=0$

$$
\begin{gathered}
\left\langle\left[\begin{array}{c}
\cos \theta \\
\sin \theta
\end{array}\right],\left[\begin{array}{c}
\cos \theta \\
\sin \theta
\end{array}\right]\right\rangle=\cos ^{2} \theta+\sin ^{2} \theta=1 \\
\left\langle\left[\begin{array}{r}
-\sin \theta \\
\cos \theta
\end{array}\right],\left[\begin{array}{r}
-\sin \theta \\
\cos \theta
\end{array}\right]\right\rangle=\cos ^{2} \theta+\sin ^{2} \theta=1 .
\end{gathered}
$$

## Example (contd.)

Let the coefficients of $\mathbf{v}$ in the first basis be $w_{1}$ and $w_{2}$ :

$$
\mathbf{v}=w_{1}\left[\begin{array}{l}
1 \\
0
\end{array}\right]+w_{2}\left[\begin{array}{l}
0 \\
1
\end{array}\right] .
$$

What are the coefficients of $\mathbf{v}$ in the second basis? Stated differently, what values of $w_{1}^{\prime}$ and $w_{2}^{\prime}$ satisfy:

$$
\mathbf{v}=w_{1}^{\prime}\left[\begin{array}{c}
\cos \theta \\
\sin \theta
\end{array}\right]+w_{2}^{\prime}\left[\begin{array}{r}
-\sin \theta \\
\cos \theta
\end{array}\right] ?
$$



Figure 1: Change of basis.
Example (contd.)
To find $w_{1}^{\prime}$ and $w_{2}^{\prime}$, we use inner product:

$$
\begin{aligned}
& w_{1}^{\prime}=\left\langle\left[\begin{array}{c}
\cos \theta \\
\sin \theta
\end{array}\right],\left[\begin{array}{l}
w_{1} \\
w_{2}
\end{array}\right]\right\rangle \\
& w_{2}^{\prime}=\left\langle\left[\begin{array}{r}
-\sin \theta \\
\cos \theta
\end{array}\right],\left[\begin{array}{l}
w_{1} \\
w_{2}
\end{array}\right]\right\rangle .
\end{aligned}
$$

## Example (contd.)

The above can be written more economically in matrix notation:

$$
\begin{gathered}
{\left[\begin{array}{l}
w_{1}^{\prime} \\
w_{2}^{\prime}
\end{array}\right]=\left[\begin{array}{rr}
\cos \theta & \sin \theta \\
-\sin \theta & \cos \theta
\end{array}\right]\left[\begin{array}{l}
w_{1} \\
w_{2}
\end{array}\right]} \\
\mathbf{w}^{\prime}=\mathbf{A w} .
\end{gathered}
$$

If the rows of $\mathbf{A}$ are orthonormal, then $\mathbf{A}$ is an orthonormal matrix. Multiplying by an orthonormal matrix effects a change of basis. A change of basis between two orthonormal bases is a rotation.

## Matrix transpose

If $\mathbf{A}$ rotates $\mathbf{w}$ by $\theta$

$$
\mathbf{A}=\left[\begin{array}{rr}
\cos \theta & \sin \theta \\
-\sin \theta & \cos \theta
\end{array}\right]
$$

then $\mathbf{A}^{-1}=\mathbf{A}^{\mathrm{T}}$ rotates $\mathbf{w}^{\prime}$ by $-\theta$

$$
\mathbf{A}^{\mathrm{T}}=\left[\begin{array}{rr}
\cos \theta & -\sin \theta \\
\sin \theta & \cos \theta
\end{array}\right]
$$

In other words, $\mathbf{A}^{\mathrm{T}}$ undoes the action of $\mathbf{A}$, i.e., they are inverses:

$$
\begin{aligned}
\mathbf{A} \mathbf{A}^{\mathrm{T}} & =\left[\begin{array}{cc}
\cos ^{2} \theta+\sin ^{2} \theta & \cos \theta \sin \theta-\sin \theta \cos \theta \\
\cos \theta \sin \theta-\sin \theta \cos \theta & \cos ^{2} \theta+\sin ^{2} \theta
\end{array}\right] \\
& =\left[\begin{array}{ll}
1 & 0 \\
0 & 1
\end{array}\right] .
\end{aligned}
$$

For orthonormal matrices, multiplying by the transpose undoes the change of basis.

## Complex vectors in $\mathbb{C}^{2}$

$\mathbf{v}=\left[a_{1} e^{i \theta_{1}}, a_{2} e^{i \theta_{2}}\right]^{\mathrm{T}}$ is a vector in $\mathbb{C}^{2}$.

- Question Can we define length and angle in $\mathbb{C}^{2}$ just like in $\mathbb{R}^{2}$ ?
- Answer Yes, but we need to redefine inner product:

$$
\langle\mathbf{u}, \mathbf{v}\rangle=u_{1}^{*} v_{1}+u_{2}^{*} v_{2} .
$$

Note that this reduces to the inner product for $\mathbb{R}^{2}$ when $\mathbf{u}$ and $\mathbf{v}$ are real. The norm of a complex vector is the square root of the sum of the squares of the amplitudes. For example, for $\mathbf{v} \in \mathbb{C}^{2}$ :

$$
\begin{aligned}
|\mathbf{u}| & =\langle\mathbf{u}, \mathbf{u}\rangle^{\frac{1}{2}} \\
& =\sqrt{u_{1}^{*} u_{1}+u_{2}^{*} u_{2}} .
\end{aligned}
$$

## Orthonormal bases for $\mathbb{C}^{n}$

- Question How about orthonormal bases for $\mathbb{C}^{n}$, do they exist?
- Answer Yes. If $\left\langle\mathbf{u}_{i}, \mathbf{u}_{j}\right\rangle=0$ when $i \neq j$ and $\left\langle\mathbf{u}_{i}, \mathbf{u}_{j}\right\rangle=1$ when $i=j$, then the $\mathbf{u}_{i}$ form an orthonormal basis for $\mathbb{C}^{n}$.
- Question Do complex orthonormal matrices exist?
- Answer Yes, except they are called unitary matrices and $\left(\mathbf{A}^{*}\right)^{\mathrm{T}}$ undoes the action of $\mathbf{A}$. That is

$$
\mathbf{A}\left(\mathbf{A}^{*}\right)^{\mathrm{T}}=\mathbf{I}
$$

where $\left(\mathbf{A}^{*}\right)^{\mathrm{T}}=\mathbf{A}^{\mathrm{H}}$ is the Hermitian transpose of $\mathbf{A}$.

## The space of $2 \pi$ periodic functions

A function, $f$, is $2 \pi$ periodic iff $f(t)=f(t+$ $2 \pi)$. We can think of two complex $2 \pi$ periodic functions, e.g., $f$ and $g$, as infinite dimensional complex vectors. Length, angle, orthogonality, and rotation (i.e., change of basis) still have meaning. All that is required is that we generalize the definition of inner product:

$$
\langle f, g\rangle=\int_{-\pi}^{\pi} f^{*}(t) g(t) d t
$$

The length (i.e., the norm) of a function is:

$$
|f|=\langle f, f\rangle^{\frac{1}{2}}=\sqrt{\int_{-\pi}^{\pi} f^{*}(t) f(t) d t}
$$

Two functions, $f$ and $g$, are orthogonal when

$$
\langle f, g\rangle=\int_{-\pi}^{\pi} f^{*}(t) g(t) d t=0
$$

## Scaling Property of the Impulse

The area of an impulse scales just like the area of a pulse, i.e., contracting an impulse by a factor of $a$ changes its area by a factor of $\frac{1}{|a|}$ :

$$
\int_{-\infty}^{\infty} \Pi(a t) d t=\frac{1}{|a|}=\lim _{\varepsilon \rightarrow 0} \int_{-\varepsilon}^{\varepsilon} \delta(a t) d t
$$

It follows that:

$$
\lim _{\varepsilon \rightarrow 0} \int_{-\varepsilon}^{\varepsilon}|a| \delta(a t) d t=\lim _{\varepsilon \rightarrow 0} \int_{-\varepsilon}^{\varepsilon} \delta(t) d t=1
$$

Since the impulse is defined by the above integral property, we conclude that:

$$
|a| \delta(a t)=\delta(t)
$$

## Shah basis

The Shah function is a train of impulses:

$$
\mathrm{III}(t)=\sum_{n=-\infty}^{\infty} \delta(t-n)
$$

We can use the scaling property of the impulse to define a $2 \pi$ periodic Shah function:

$$
\begin{aligned}
\frac{1}{2 \pi} \mathrm{III}\left(\frac{t}{2 \pi}\right) & =\frac{1}{2 \pi} \sum_{n=-\infty}^{\infty} \delta\left(\frac{t}{2 \pi}-n\right) \\
& =\frac{2 \pi}{2 \pi} \sum_{n=-\infty}^{\infty} \delta\left(2 \pi\left(\frac{t}{2 \pi}-n\right)\right) \\
& =\sum_{n=-\infty}^{\infty} \delta(t-2 \pi n) .
\end{aligned}
$$

## Shah basis (contd.)

Consider the infinite set of $2 \pi$ periodic Shah functions, $\frac{1}{2 \pi} \operatorname{III}\left(\frac{t-\tau}{2 \pi}\right)$, for $-\pi \leq \tau<\pi$. Because $\frac{1}{2 \pi} \operatorname{III}\left(\frac{t-\tau}{2 \pi}\right)=\delta(t-\tau)$ for $-\pi \leq t \leq \pi$ it follows that

$$
\begin{gathered}
\left\langle\frac{1}{2 \pi} \mathrm{III}\left(\frac{t-\tau_{1}}{2 \pi}\right), \frac{1}{2 \pi} \operatorname{III}\left(\frac{t-\tau_{2}}{2 \pi}\right)\right\rangle \\
=\int_{-\pi}^{\pi} \delta\left(t-\tau_{1}\right) \delta\left(t-\tau_{2}\right) d t
\end{gathered}
$$

equals 0 when $\tau_{1} \neq \tau_{2}$ and equals $\int_{-\pi}^{\pi} \delta\left(t-\tau_{1}\right) d t=$ 1 when $\tau_{1}=\tau_{2}$. It follows that the infinite set of $2 \pi$ periodic Shah functions, $\frac{1}{2 \pi} \operatorname{III}\left(\frac{t-\tau}{2 \pi}\right)$, for $-\pi \leq \tau<\pi$ form an orthonormal basis for the space of $2 \pi$ periodic functions.
$\operatorname{III}(t)=\sum_{n=-\infty}^{\infty} \delta(t-n)$


$$
\frac{1}{2 \pi} \operatorname{III}\left(\frac{t-\tau}{2 \pi}\right)=\sum_{n=-\infty}^{\infty} \delta(t-\tau-n 2 \pi)
$$



Figure 2: Making a $2 \pi$ periodic Shah function.

## Shah basis (contd.)

- Question How do we find the coefficients, $w(\tau)$, representing $f(t)$ in the Shah basis? How do we find $w(\tau)$ such that

$$
f(t)=\frac{1}{2 \pi} \int_{-\pi}^{\pi} w(\tau) \mathrm{III}\left(\frac{t-\tau}{2 \pi}\right) d \tau ?
$$

- Answer Take inner products of $f$ with the infinite set of $2 \pi$ periodic Shah functions:

$$
w(\tau)=\left\langle\frac{1}{2 \pi} \operatorname{III}\left(\frac{t-\tau}{2 \pi}\right), f(t)\right\rangle .
$$

## Shah basis (contd.)

Because $\frac{1}{2 \pi} \operatorname{III}\left(\frac{t-\tau}{2 \pi}\right)=\delta(t-\tau)$ for $-\pi \leq t \leq \pi$ it follows that

$$
\begin{aligned}
w(\tau) & =\frac{1}{2 \pi} \int_{-\pi}^{\pi} f(t) \operatorname{III}\left(\frac{t-\tau}{2 \pi}\right) d t \\
& =\int_{-\pi}^{\pi} f(t) \delta(t-\tau) d t
\end{aligned}
$$

which by the sifting property of the impulse is just:

$$
w(\tau)=f(\tau)
$$

We see that the coefficients of $f$ in the Shah basis are just $f$ itself!

## Harmonic signal basis

- Question How long is a harmonic signal?
- Answer The length of a harmonic signal is

$$
\begin{aligned}
\left|e^{j \omega t}\right| & =\left\langle e^{j \omega t}, e^{j \omega t}\right\rangle^{\frac{1}{2}} \\
& =\left(\int_{-\pi}^{\pi} e^{-j \omega t} e^{j \omega t} d t\right)^{\frac{1}{2}} \\
& =\left(\int_{-\pi}^{\pi} d t\right)^{\frac{1}{2}} \\
& =\sqrt{2 \pi} .
\end{aligned}
$$

Harmonic signal basis (contd.)

- Question What is the angle between two harmonic signals with integer frequencies?
- Answer The angle between two harmonic signals with integer frequencies is

$$
\begin{aligned}
\left\langle e^{j \omega_{1} t}, e^{j \omega_{2} t}\right\rangle & =\int_{-\pi}^{\pi} e^{-j \omega_{1} t} e^{j \omega_{2} t} d t \\
& =\left.\left[\frac{e^{j\left(\omega_{2}-\omega_{1}\right) t}}{j\left(\omega_{2}-\omega_{1}\right)}\right]\right|_{-\pi} ^{\pi}
\end{aligned}
$$

Since this function is the same at $-\pi$ and $\pi$ (for all integers $\omega_{1}$ and $\omega_{2}$ ), we conclude that

$$
\left\langle e^{j \omega_{1} t}, e^{j \omega_{2} t}\right\rangle=0
$$

when $\omega_{1}$ and $\omega_{2}$ are integers and $\omega_{1} \neq \omega_{2}$.

## Fourier Series of $2 \pi$ Periodic Functions

It follows that the infinite set of harmonic signals, $\frac{1}{\sqrt{2 \pi}} e^{j \omega t}$ for integer $\omega$ and $-\infty \leq \omega \leq \infty$ form an orthonormal basis for the space of $2 \pi$ periodic functions.

- Question What are the coefficients of $f$ in the harmonic signal basis?
- Answer Take inner products of $f$ with the infinite set of harmonic signals.

This is the analysis formula for Fourier series:

$$
\begin{aligned}
F(\omega) & =\left\langle\frac{1}{\sqrt{2 \pi}} e^{j \omega t}, f\right\rangle \\
& =\frac{1}{\sqrt{2 \pi}} \int_{-\pi}^{\pi} f(t) e^{-j \omega t} d t
\end{aligned}
$$

for integer frequency, $\omega$.

## Fourier Series of $2 \pi$ Periodic Functions (contd.)

The function can be reconstructed using the synthesis formula for Fourier series:

$$
f(t)=\frac{1}{\sqrt{2 \pi}} \sum_{\omega=-\infty}^{\infty} F(\omega) e^{j \omega t}
$$

## Fourier Series Example

The Fourier series for the Shah basis function

$$
f(t)=\frac{1}{2 \pi} \mathrm{III}\left(\frac{t}{2 \pi}\right)
$$

is

$$
\begin{aligned}
F(\omega) & =\left\langle\frac{1}{\sqrt{2 \pi}} e^{j \omega t}, \frac{1}{2 \pi} \operatorname{III}\left(\frac{t}{2 \pi}\right)\right\rangle \\
& =\frac{1}{\sqrt{2 \pi}} \int_{-\pi}^{\pi} \delta(t) e^{j \omega t} d t \\
& =\frac{1}{\sqrt{2 \pi}}
\end{aligned}
$$

Consequently

$$
\begin{aligned}
f(t) & =\frac{1}{\sqrt{2 \pi}} \sum_{\omega=-\infty}^{\infty} F(\omega) e^{j \omega t} \\
& =\frac{1}{2 \pi} \sum_{\omega=-\infty}^{\infty} e^{j \omega t}
\end{aligned}
$$

## Deep Thought

The analysis formula for Fourier series effects a change of basis. It is a rotation in the space of $2 \pi$ periodic functions. The synthesis formula undoes the change of basis. It is the opposite rotation.


Figure 3: $\operatorname{Re}(t)$ (solid) and $\operatorname{Im}(t)$ (dashed) of truncated Fourier series for Shah basis function. (a) $-1 \leq \omega \leq 1$ (b) $-2 \leq \omega \leq 2$ (c) $-4 \leq \omega \leq 4$ (d) $-8 \leq \omega \leq 8$ (e) $-16 \leq \omega \leq 16$ (f) $-32 \leq \omega \leq 32$.


Figure 4: $\operatorname{Re}(t)$ (solid) and $\operatorname{Im}(t)$ (dashed) of truncated Fourier series for Shah basis function. (a) $-64 \leq \omega \leq 64$ (b) $-128 \leq \omega \leq 128$ (c) $-256 \leq \omega \leq 256$ (d) $-512 \leq \omega \leq 512$ (e) $-1024 \leq \omega \leq$ 1024 (f) $-2048 \leq \omega \leq 2048$.

## Fourier Series of $T$-Periodic Functions

A function, $f$, is $T$-periodic iff $f(t)=f(t+T)$.

- Analysis formula

$$
\begin{aligned}
F(\omega) & =\left\langle\frac{\sqrt{2 \pi}}{T} e^{j 2 \pi \omega t / T}, f\right\rangle \\
& =\frac{\sqrt{2 \pi}}{T} \int_{-T / 2}^{T / 2} f(t) e^{-j 2 \pi \omega t / T} d t
\end{aligned}
$$

for integer frequency, $\omega$.

- Synthesis formula

$$
f(t)=\frac{\sqrt{2 \pi}}{T} \sum_{\omega=-\infty}^{\infty} F(\omega) e^{j 2 \pi \omega t / T}
$$

Observe that if we substitute $T=2 \pi$ in the above expressions, we get the formulas for $2 \pi$ periodic functions.

## The Fourier Transform

Functions with finite length are termed square integrable.

$$
\begin{aligned}
|f| & =\sqrt{\int_{-\infty}^{\infty}|f(t)|^{2} d t} \\
& =\sqrt{\int_{-\infty}^{\infty} f^{*}(t) f(t) d t} \\
& <\infty .
\end{aligned}
$$

For square integrable functions, we can take the limit of the Fourier series for $T$-periodic functions as $T \rightarrow \infty$, in which case, it is possible to show that...

## The Fourier Transform (contd.)

- Analysis formula

$$
\begin{aligned}
F(s) & =\left\langle e^{j 2 \pi s t}, f\right\rangle \\
& =\int_{-\infty}^{\infty} f(t) e^{-j 2 \pi s t} d t
\end{aligned}
$$

- Synthesis formula

$$
f(t)=\int_{-\infty}^{\infty} F(s) e^{j 2 \pi s t} d s
$$

